LECTURE 1 INTRODUCTION TO QUANTUM COMPUTING

Quantum computer science and applications

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ORIGIN OF QUANTUM COMPUTING

Feynman (1981):

Can quantum systems be probabilistically simulated by a classical computer?

→ The answer is almost certainly, no!

 \longrightarrow Use quantum systems/computers to simulate quantum systems!

(birth of quantum simulation)

A natural question:

What other problems can quantum computers solve more quickly than classical computer?

Deutsch (1985):

Foundation of quantum computing!

→ Deutsch-Jozsa algorithm (1992) quantum algorithm faster than any classical algorithm

EARLY ALGORITHMS: SHOR

Shor (1994):

Solves the discrete logarithm and factoring problem efficiently with a quantum computer!

Terrible situation: public-key cryptography currently deployed is broken by using an "efficient" quantum computer

 \longrightarrow Cryptographic community worried about this since many years. . .

There exists quantum resistant solutions: post-quantum cryptography (active research topic)

American's government (2017 & 2023) has launched processes to standardized post-quantum cryptosystems

EARLY ALGORITHMS: GROVER

Grover (1996):

Find an element in a list of size n in time $O\left(\sqrt{n}\right)$ while any classical algorithm needs a time $\approx n$

Consequence: size of keys in symmetric cryptography has to be ×2.

 $\left(\text{ size of cryptosystem }\ell\text{ bits: best classical attack costs }2^{\ell}\stackrel{\text{(Grover)}}{\longmapsto}2^{\ell/2}\right)$

A BIG ISSUE: DECOHERENCE

Computations are "noisy"

- Quantum bits are very fragile, they quickly interfere with the environment: decoherence
- Quantum architectures are not "ideal"
 - \longrightarrow Faults in computation can theoretically be "corrected": quantum error correcting codes

Theorem [Aharonov, Ben-Or, 1997]:

Quantum computation is possible provided the noise is sufficiently low

QUANTUM CRYPTOGRAPHY

Benett-Brassard (1984):

Quantum protocol for key-exchange

- ► Already implemented
- ► If an authenticated canal has been established, unconditional security: relies strongly on the validity of physic laws and not computational assumptions

PROGRAM OF THIS COURSE

→ Basics of quantum computing and quantum information theory

- Quantum formalism with density operators, general measures, partial trace, etc. . .
- Quantum circuit model, quantum algorithms (Deutsch-Josza, Simon, Grover, Quantum Fourier Transform, Shor, Kitaev)
- Basics of quantum error correcting codes and quantum cryptography

References:

- ▶ Nielsen and Chuang, Quantum computation and quantum information,
 - → Nice introduction to quantum computing and quantum information
- ▶ de Wolf's lecture notes: https://arxiv.org/abs/1907.09415,
 - \longrightarrow Nice for advanced quantum algorithms
- ► Childs's lecture notes: https://www.cs.umd.edu/~amchilds/qa/,
 - → Nice for advanced topics
- ► Zemor's lecture notes: https://www.math.u-bordeaux.fr/~gzemor/QuantumCodes.pdf,
 - → Introduction to quantum error correcting codes

EVALUATION OF THIS COURSE

- 1. An exam (3 hours): 4 pages of personal notes are allowed
 - \longrightarrow Three exercises seen during the Exercise Sessions will be at the exam
- 2. Presentation of a research article or a chapter of some lecture notes (30min)

You are in a course of computer science

Computer science: art of computing

→ We don't care that an object "exists", we want to compute it efficiently!

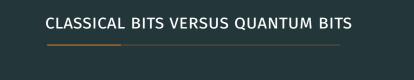
Using the law of quantum physic: new model of computation

What does mean quantum computing? What is a quantum algorithm?

 \longrightarrow This course is not about the law of physics or about the "technologies" to verify/use them

COURSE OUTLINE

- 1. Classical Bits Versus Quantum Bits
- 2. Your First Quantum Algorithm
- 3. n-qubits Systems
- 4. Bra-ket and Ket-bra Notation



CLASSICAL BIT

- ► Classical bit: $b \in \{0,1\}$ with XOR operation $(1 \oplus 1 = 0 \oplus 0 = 0 \text{ and } 1 \oplus 0 = 0 \oplus 1 = 1)$
- Probabilistic bit: $\binom{p}{q}$ where

$$p \stackrel{\text{def}}{=} \mathbb{P}(b=0)$$
$$q \stackrel{\text{def}}{=} \mathbb{P}(b=1)$$

► Evolution during a computation (a probabilistic bit stays a probabilistic bit):

$$\begin{pmatrix} p \\ q \end{pmatrix} \longrightarrow \begin{pmatrix} p' \\ q' \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} p \\ q \end{pmatrix} \quad \text{where } \left\{ \begin{array}{l} a+c=1 \\ b+d=1 \end{array} \right. \text{ and } a,b,c,d \geq 0.$$

Probabilistic computation: multiplication by a stochastic matrix

Examples: $b \rightarrow b \oplus b$ and $b \mapsto b \oplus 1$

$$\begin{pmatrix} p \\ q \end{pmatrix} \longrightarrow \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} p \\ q \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} p \\ q \end{pmatrix} \longrightarrow \begin{pmatrix} q \\ p \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} p \\ q \end{pmatrix}$$

"A superposition of classical states"

▶ A qubit $|\psi\rangle$ is an element of \mathbb{C}^2 with Hermitian norm 1:

$$|\psi\rangle=\alpha\,|0\rangle+\beta\,|1\rangle$$
 with $\alpha,\beta\in\mathbb{C}$ (called amplitude) and $|\alpha|^2+|\beta|^2=1$

where $(|0\rangle, |1\rangle)$ orthonormal basis of \mathbb{C}^2 . Usually defined as

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$
 and $|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

"A superposition of classical states"

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$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$
 and $|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

We "cannot see" a superposition, we "can only see" classical states: measure and observe!

▶ Measurement: probabilistic orthogonal projection. Given $|e_0\rangle$, $|e_1\rangle \in \mathbb{C}^2$ orthonormal basis:

Measuring in the basis
$$(|e_0\rangle\,,|e_1\rangle): |\psi\rangle = \alpha\,|e_0\rangle\,+\beta\,|e_1\rangle \xrightarrow{measure} \left\{ \begin{array}{l} |e_0\rangle \text{ with prob. } |\alpha|^2 \\ |e_1\rangle \text{ with prob. } |\beta|^2 \end{array} \right.$$

Exercise: Computational versus Hadamard basis

1. Show that ($|+\rangle$, $|-\rangle$) is an orthonormal basis of \mathbb{C}^2 where

$$|+\rangle\stackrel{\text{def}}{=}\frac{1}{\sqrt{2}}\left(|0\rangle+|1\rangle\right)$$
 and $|-\rangle\stackrel{\text{def}}{=}\frac{1}{\sqrt{2}}\left(|0\rangle-|1\rangle\right)$

2. Give the outcome distribution when measuring $|0\rangle$, $|-\rangle$, and $\frac{1}{\sqrt{3}}|0\rangle + \sqrt{\frac{2}{3}}|1\rangle$ in the bases $(|0\rangle, |1\rangle)$ and $(|+\rangle, |-\rangle)$

- ▶ Qubit: $|\psi\rangle \in \mathbb{C}^2$ of Hermitian norm 1
- ▶ Measuring in the orthonormal basis ($|e_0\rangle$, $|e_1\rangle$):

$$|\psi\rangle = \alpha\,|e_0\rangle + \beta\,|e_1\rangle \xrightarrow{\text{measure}} \left\{ \begin{array}{l} |e_0\rangle \ \ \text{with probability} \ |\alpha|^2 \\ |e_1\rangle \ \ \text{with probability} \ |\beta|^2 \end{array} \right.$$

A measurement is a computation you have access to

 \longrightarrow See Lecture 2 for a precise definition of measurement. . .

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Are there other computations over qubits we have access to?

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Are there other computations over qubits we have access to?

→ Yes! Unitary evolutions

UNITARY EVOLUTIONS

$$\mathbf{U} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \mathbb{C}^{2 \times 2}, \text{ then its conjugate transpose } \mathbf{U}^\dagger = \begin{pmatrix} \overline{a} & \overline{c} \\ \overline{b} & \overline{d} \end{pmatrix}$$

▶ Unitary evolution: $U \in \mathbb{C}^{2 \times 2}$ unitary matrix $\iff UU^{\dagger} = I_2$

$$|\psi\rangle \longrightarrow \mathsf{U}\,|\psi\rangle$$

Is it true that a qubit is still a qubit after a unitary evolution? Why?

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Is it true that a qubit is still a qubit after a unitary evolution? Why?

 \longrightarrow Yes! Unitary evolutions preserve the Hermitian norm (more generally the Hermitian product)

Unitary evolutions are invertible!

$$|\psi\rangle \stackrel{\mathsf{U}}{\longrightarrow} \mathsf{U}\,|\psi\rangle \stackrel{\mathsf{U}^{\dagger}}{\longrightarrow} \mathsf{U}^{\dagger} \mathsf{U}\,|\psi\rangle = |\psi\rangle$$

- ▶ $U \in \mathbb{C}^{2 \times 2}$ unitary over qubits is often called quantum gate
 - \longrightarrow It exists a small set of gates which is universal (be patient, wait Lecture 4)

To define a quantum gate: enough to specify the image of an orthonormal basis and then extended it by linearity

But it has to map an orthonormal basis to an orthonormal basis!

Exercise: Quantum Gates?

Are the following linear operators over qubits be quantum gates?

1.
$$|0\rangle \mapsto |1\rangle$$
 and $|1\rangle \mapsto \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$

2.
$$|0\rangle \mapsto |1\rangle$$
 and $|1\rangle \mapsto |0\rangle$

Quantum gates have matrix representations!

For instance: $|0\rangle \mapsto |1\rangle$ and $|1\rangle \mapsto |0\rangle$ has the representation: $\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$. Only linear operator that maps $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ to $|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ and $|1\rangle$ to $|0\rangle$

EXAMPLE OF QUANTUM GATES

► NOT-gate X:

Linear op.	Matrix rep.
0⟩ → 1⟩	$\begin{pmatrix} 0 & 1 \end{pmatrix}$
1⟩ → 0⟩	1 0

► Hadamard-gate H:

Linear op.	Matrix rep.
$ 0\rangle \mapsto \frac{1}{\sqrt{2}} (0\rangle + 1\rangle)$ $ 1\rangle \mapsto \frac{1}{\sqrt{2}} (0\rangle - 1\rangle)$	$ \begin{array}{c c} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \end{array} $

Exercise:

- 1. What is the effect of applying \boldsymbol{H} on $|0\rangle$ and measuring it?
- 2. What is the effect of applying **H** on $|0\rangle$ twice?

CLASSICAL VERSUS QUANTUM COMPUTATION

Is quantum computation over qubits the same than classical computation over probabilistic bits?

Exercise:

Show that there is no stochastic matrix **P** which when applied to 0, *i.e.* to $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$, simulates the effect of the Hadamard gate

The "-1" gives you a huge power. . .



THE DEUTSCH-JOSZA PROBLEM

Problem:

- Input: $f: \{0,1\}^n \to \{0,1\}$ either constant or balanced
- Output: 0 if and only if f is constant

Query complexity to f to find the correct answer with certainty:

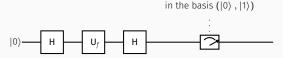
- ► Classically: $1 + \frac{2^n}{2}$
- Quantumly: 1

THE DEUTSCH-JOSZA ALGORITHM FOR n=1

Suppose that we have access to the following gate (see Exercise Session)

$$|b\rangle$$
 $(-1)^{f(b)}|b\rangle$

► The algorithm



- Analysis
- 1. Applying H: $\frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$
- 2. Applying **U**_f:

$$U_{f}\left(\frac{1}{\sqrt{2}}\left(\left|0\right\rangle + \left|1\right\rangle\right)\right) = \frac{1}{\sqrt{2}}\left(U_{f}\left|0\right\rangle + U_{f}\left|1\right\rangle\right) = \frac{(-1)^{f(0)}\left|0\right\rangle + (-1)^{f(1)}\left|1\right\rangle}{\sqrt{2}}$$

3. Applying H:

$$H\left(\frac{(-1)^{f(0)}|0\rangle + (-1)^{f(1)}|1\rangle}{\sqrt{2}}\right) = \frac{1}{\sqrt{2}}\left((-1)^{f(0)}H|0\rangle + (-1)^{f(1)}H|1\rangle\right)$$
$$= \frac{\left((-1)^{f(0)} + (-1)^{f(0)}\right)|0\rangle + \left((-1)^{f(0)} - (-1)^{f(1)}\right)|1\rangle}{2}$$

Before measuring we have computed:

$$|\psi_{\text{out}}\rangle \stackrel{\text{def}}{=} \frac{\left((-1)^{f(0)} + (-1)^{f(1)}\right)|0\rangle + \left((-1)^{f(0)} - (-1)^{f(1)}\right)|1\rangle}{2}$$

► If *f* constant:

$$|\psi_{
m out}
angle=\pm\,|0
angle$$

▶ If f balanced, namely $f(0) \neq f(1)$:

$$|\psi_{
m out}
angle=\pm\,|1
angle$$

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m out}
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▶ If f balanced, namely $f(0) \neq f(1)$:

$$|\psi_{
m out}
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Measuring in the ($|0\rangle$, $|1\rangle$) basis leads to (with probability one)

 $|0\rangle$ if f constant or $|1\rangle$ if f balanced

n **QUBITS SYSTEMS**

FINITE DIMENSION

During all this course we will work in finite dimension, think \mathbb{C}^N

 \longrightarrow Vector spaces have finite dimension, linear operators can be written as matrices, etc. . .

TENSOR PRODUCT

Given two vector spaces V and W, the tensor product $\mathbf{v} \otimes \mathbf{w}$ between $\mathbf{v} \in V$ and $\mathbf{w} \in W$ verifies:

(1) for any scalar z,

$$z(v \otimes w) = (zv) \otimes w = v \otimes (zw)$$

(2) for any $\mathbf{v}_1, \mathbf{v}_2 \in V$,

$$(v_1 + v_2) \otimes w = v_1 \otimes w + v_2 \otimes w$$

(3) for any $\mathbf{w}_1, \mathbf{w}_2 \in W$,

$$v\otimes (w_1+w_2)=v\otimes w_1+v\otimes w_2$$

The tensor product $\mathbf{v} \otimes \mathbf{w}$ as a column/row product:

$$\begin{pmatrix} v_1 \\ \vdots \\ v_N \end{pmatrix} \begin{pmatrix} w_1 & \cdots & w_N \end{pmatrix}$$

TENSOR PRODUCT OF SPACES

Tensor product of spaces:

V and W be two vector spaces with bases the \mathbf{v}_i 's and the \mathbf{w}_j respectively

$$V = \text{Span}(\mathbf{v}_1, \dots, \mathbf{v}_n)$$
 and $W = \text{Span}(\mathbf{w}_1, \dots, \mathbf{w}_m)$

The vector space $V \otimes W$ is defined as being generated by the \mathbf{v}_i 's and the \mathbf{w}_j 's

$$V \otimes W \stackrel{\text{def}}{=} \text{Span} (\mathbf{v}_i \otimes \mathbf{w}_j : 1 \leq i \leq n, 1 \leq j \leq m)$$

► Dimension is multiplicative

$$\dim V \otimes W = \dim V \dim W = nm$$

Basis, $(\mathbf{v}_1, \dots, \mathbf{v}_n)$ (resp. $(\mathbf{w}_1, \dots, \mathbf{w}_m)$) be a basis of V (resp. W)

$$(\mathbf{v}_i \otimes \mathbf{w}_j : 1 \le i \le n, \ 1 \le j \le m)$$
 is a basis of $V \otimes W$

Characterization

$$\mathbf{x} \in V \otimes W \iff \exists \alpha_{i,j} \ : \ \mathbf{x} = \sum_{\substack{1 \leq i \leq n \\ 1 \leq i \leq m}} \alpha_{i,j} \ \mathbf{v}_i \otimes \mathbf{w}_j$$

Classical error:

 $\mathbf{x} \in V \otimes W$, then it exists $\mathbf{v} \in V$ and $\mathbf{w} \in W$ such that $\mathbf{x} = \mathbf{v} \otimes \mathbf{w}$.

SCALAR PRODUCT OVER TENSOR PRODUCT SPACES

$$(\mathbf{v}_1,\ldots,\mathbf{v}_n)$$
 $\Big($ resp. $(\mathbf{w}_1,\ldots,\mathbf{w}_m)\Big)$ be a basis of V $\Big($ resp. W $\Big)$

Scalar product over tensor product spaces:

Suppose that V (resp. W) is equipped by a scalar product $\langle \cdot, \cdot \rangle_V$ (resp. $\langle \cdot, \cdot \rangle_W$). The scalar product over $V \otimes W$ is defined as (and extended by bilinearity) as:

$$\langle \mathbf{v}_i \otimes \mathbf{w}_j, \mathbf{v}_k \otimes \mathbf{w}_\ell \rangle_{V \otimes W} \stackrel{\mathsf{def}}{=} \langle \mathbf{v}_i, \mathbf{v}_k \rangle_V \cdot \langle \mathbf{w}_j, \mathbf{w}_\ell \rangle_W$$

An important remark:

If
$$\mathbf{v}_1 \perp \mathbf{v}_2$$
, then for all $\mathbf{w}_1, \mathbf{w}_2$: $(\mathbf{v}_1 \otimes \mathbf{w}_1) \perp (\mathbf{v}_2 \otimes \mathbf{w}_2)$

LINEAR OPERATOR OVER TENSOR PRODUCT OF SPACES

$$(\textbf{v}_1,\ldots,\textbf{v}_n)$$
 (resp. $(\textbf{w}_1,\ldots,\textbf{w}_m)$ be a basis of V (resp. W)

Linear operator over tensor product of spaces:

Given A, B be linear operators over V, W, $A \otimes B$ is a linear operator over $V \otimes W$ be defined (and extended by linearity) as:

$$A \otimes B \left(v_i \otimes w_j \right) \stackrel{\text{def}}{=} A v_i \otimes B w_j$$

Characterization:

C linear operator over
$$V \otimes W \iff \exists \alpha_i, \mathsf{A}_i, \mathsf{B}_i \; : \; \mathsf{C} = \sum_i \alpha_i \; \mathsf{A}_i \otimes \mathsf{B}_i$$

Classical error:

C linear operator over V \otimes W, then there exists A, B linear operators over V and W s.t $C=A\otimes B$

Tensor product of matrices:

Let $\mathbf{A} \stackrel{\text{def}}{=} (a_{i,j})_{\substack{1 \leq i \leq n \\ 1 \leq j \leq m}} \in \mathbb{C}^{n \times m}$ and $\mathbf{B} \in \mathbb{C}^{p \times q}$, then

$$\mathbf{A} \otimes \mathbf{B} \stackrel{\text{def}}{=} \begin{pmatrix} a_{1,1}\mathbf{B} & a_{1,2}\mathbf{B} & \cdots & a_{1,m}\mathbf{B} \\ a_{2,1}\mathbf{B} & a_{2,2}\mathbf{B} & \cdots & a_{2,m}\mathbf{B} \\ \vdots & \vdots & \vdots & \vdots \\ a_{n,1}\mathbf{B} & a_{n,2}\mathbf{B} & \cdots & a_{n,m}\mathbf{B} \end{pmatrix} \in \mathbb{C}^{np \times mq}$$

Example:

1.
$$\binom{1}{2} \otimes \binom{2}{3} = \binom{1 \times 2}{1 \times 3}_{2 \times 2}_{2 \times 3} = \binom{2}{3}_{4}_{6}.$$

2.
$$\mathbf{X} \otimes \mathbf{H} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \otimes \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \\ 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \end{pmatrix}$$

PROPERTIES OF THE TENSOR PRODUCT OF MATRICES

Properties:

For any $\alpha \in \mathbb{C}$, $A, B \in \mathbb{C}^{m \times n}$ and $C, D \in \mathbb{C}^{p \times q}$

- 1. $\alpha (A \otimes C) = (\alpha A) \otimes C = A \otimes (\alpha C)$
- 2. $(A + B) \otimes C = A \otimes C + B \otimes C$
- 3. $C \otimes (A + B) = C \otimes A + C \otimes B$
- 4. If we can form matrices products AC and BD, then

$$(A \otimes B) (C \otimes D) = (AC) \otimes (BD)$$

5. If A, B are invertible, then

$$(\mathbf{A} \otimes \mathbf{B})^{-1} = \mathbf{A}^{-1} \otimes \mathbf{B}^{-1}.$$

Classical error:

$$A \otimes B = B \otimes A$$

n QUBITS SYSTEMS

- lacktriangle A qubit $|\psi\rangle$ is an element of \mathbb{C}^2 with Hermitian norm 1
- A register of *n* qubits $|\psi\rangle$ is an element of $\underbrace{\mathbb{C}^2 \otimes \cdots \otimes \mathbb{C}^2}_{n \text{ times}} = \mathbb{C}^{2^n}$ with Hermitian norm 1

Let $(|0\rangle, |1\rangle)$ be an orthonormal basis of \mathbb{C}^2 . Then,

$$(|b_1\rangle \otimes |b_2\rangle \otimes \cdots \otimes |b_n\rangle : b_1, \ldots, b_n \in \{0,1\})$$

is an orthonormal basis of $\underbrace{\mathbb{C}^2 \otimes \cdots \otimes \mathbb{C}^2}_{n \text{ times}} = \mathbb{C}^{2^n}$

Notation: for $\mathbf{b} = (b_1, \dots, b_n) \in \{0, 1\}^n$ and $|\psi_1\rangle, |\psi_2\rangle, \dots, |\psi_n\rangle$ be qubits

$$\mathbf{b} = |b_1 b_2 \dots b_n\rangle \stackrel{\text{def}}{=} |b_1\rangle \otimes |b_2\rangle \otimes \dots \otimes |b_n\rangle \quad \text{and} \quad |\psi_1\rangle |\psi_2\rangle \dots |\psi_n\rangle \stackrel{\text{def}}{=} |\psi_1\rangle \otimes |\psi_2\rangle \otimes \dots \otimes |\psi_n\rangle$$

• Characterization: any register $|\psi\rangle\in\mathbb{C}^{2^n}$ of n qubits can be written as

$$|\psi\rangle = \sum_{\mathbf{x} \in \{0,1\}^n} \alpha_{\mathbf{x}} \, |\mathbf{x}\rangle \quad \text{ where } \alpha_{\mathbf{x}} \in \mathbb{C} \, \left(\text{called amplitude} \right) \quad \text{and } \sum_{\mathbf{x} \in \{0,1\}^n} |\alpha_{\mathbf{x}}|^2 = 1$$

n QUBITS SYSTEMS

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A remark: choose your orthonormal basis!

From any (
$$|e_0\rangle$$
, $|e_1\rangle$) orthonormal basis of \mathbb{C}^2 , then $\left(\left|e_{i_1}\right\rangle\otimes\cdots\otimes\left|e_{i_n}\right\rangle\right)$ for $i_1,\ldots,i_n\in\{0,1\}^n$ is an orthonormal basis of $\mathbb{C}^2\otimes\cdots\otimes\mathbb{C}^2=\mathbb{C}^{2^n}$

Exercise:

- 1. Compute the scalar product between $|+\rangle |1\rangle$, $|00\rangle$ and $|11\rangle$ where $|+\rangle = \frac{1}{\sqrt{2}} (|0\rangle |1\rangle)$
- 2. Let $(|e_0\rangle, |e_1\rangle)$ be an orthonormal basis of \mathbb{C}^2 . Show that $(|e_{i_1}\rangle \dots |e_{i_n}\rangle)$ for $i_1, \dots, i_n \in \{0, 1\}^n$ is an orthonormal basis of $\underbrace{\mathbb{C}^2 \otimes \dots \otimes \mathbb{C}^2}_{n \text{ times}} = \mathbb{C}^{2^n}$
- 3. Do we have $|00\rangle + |10\rangle = (|0\rangle + |1\rangle) \otimes |0\rangle$?
- 4. (*) Do there exist two qubits $|\psi_1\rangle$ and $|\psi_2\rangle$ such that

$$\frac{1}{\sqrt{2}}\left(|00\rangle+|11\rangle\right)=|\psi_1\rangle\otimes|\psi_2\rangle$$

5. Do there exist two qubits $|\psi_1\rangle$ and $|\psi_2\rangle$ such that

$$\frac{1}{2}\left(|00\rangle+|01\rangle+|10\rangle+|11\rangle\right)=|\psi_1\rangle\otimes|\psi_2\rangle$$

Separable versus entangled states:

A n-qubit system $|\psi\rangle$ that can be decomposed as $|\psi\rangle = |\psi_1\rangle \otimes |\psi_2\rangle$ is called separable When there is no such decomposition, the state is called entangled

Example:

1. Separable states,

$$|00\rangle = |0\rangle \otimes |0\rangle \quad \text{and} \quad \frac{1}{2} \left(|00\rangle + |01\rangle + |10\rangle + |11\rangle \right) = \frac{1}{\sqrt{2}} \left(|0\rangle + |1\rangle \right) \otimes \frac{1}{\sqrt{2}} \left(|0\rangle + |1\rangle \right)$$

2. Entangled state,

$$\frac{1}{\sqrt{2}}(|00\rangle+|11\rangle)$$

→ Entangled states play a crucial role in quantum computation/information (teleportation,

quantum cryptography, . . .)

► Measuring in the basis $|e_1\rangle |e_2\rangle \cdots |e_n\rangle$:

$$\left|\psi\right\rangle = \sum_{i_1, \dots, i_n \in \left\{0,1\right\}^n} \alpha_{i_1 \dots i_n} \left|e_{i_1}\right\rangle \cdots \left|e_{i_n}\right\rangle \xrightarrow{\textit{measure}} \left|e_{j_1}\right\rangle \cdots \left|e_{j_n}\right\rangle \text{ with probability } \left|\alpha_{j_1 \dots j_n}\right|^2$$

► Measuring the first register in the basis ($|e_0\rangle$, $|e_1\rangle$):

$$\left|\psi\right\rangle = \alpha_0 \left|e_0\right\rangle \left|\psi_0\right\rangle + \alpha_1 \left|e_1\right\rangle \left|\psi_1\right\rangle \xrightarrow{\textit{measure}} \left\{ \begin{array}{c} \left|e_0\right\rangle \left|\psi_0\right\rangle & \textit{with probability } \left|\alpha_0\right|^2 \\ \left|e_1\right\rangle \left|\psi_1\right\rangle & \textit{with probability } \left|\alpha_1\right|^2 \end{array} \right.$$

Be careful: necessarily
$$|\alpha_0|^2 + |\alpha_1|^2 = 1$$
.

Exercise:

Give the outcome distribution of measuring in the basis ($|bb'\rangle:b,b'\in\{0,1\}$) the first registers of the following two-qubits:

$$|0\rangle \left(\sqrt{\frac{1}{3}}|0\rangle + \sqrt{\frac{2}{3}}|1\rangle\right), \quad \sqrt{\frac{1}{2}}|01\rangle + \sqrt{\frac{1}{3}}|11\rangle + \sqrt{\frac{1}{6}}|10\rangle \quad \text{and} \quad \frac{1}{2}(|0\rangle - |1\rangle)(|0\rangle - |1\rangle)$$

UNITARY FOR *n* QUBITS SYSTEMS

Unitary evolution
$$\mathbf{U} \in \mathbb{C}^{2^n \times 2^n}$$
 unitary matrix $\iff \mathbf{U}\mathbf{U}^\dagger = \mathbf{Id}_{2^n}$

Exercise:

Is the following operator a unitary of $\mathbb{C}^2\otimes\mathbb{C}^2$:

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

Describe the image of $\left|bb'\right>$ for $b,b'\in\{0,1\}$

BRA-KET AND KET-BRA NOTATION

THE BRA-KET NOTATION

Scalar Product:

Let $|e_1\rangle$, ..., $|e_{2^n}\rangle$ be an orthonormal basis, $|\psi\rangle\stackrel{\mathrm{def}}{=} \sum_i \alpha_i \, |e_i\rangle$ and $|\varphi\rangle\stackrel{\mathrm{def}}{=} \sum_i \beta_i \, |e_i\rangle$. Then $\langle\psi|\varphi\rangle\stackrel{\mathrm{def}}{=} \sum_i \overline{\alpha_i}\beta_i$

- Ket-notation: $|\psi\rangle$ is called a ket
- ▶ Bra-notation: a ket $|\psi\rangle = \begin{pmatrix} \alpha_1 \\ \vdots \\ \alpha_{2^n} \end{pmatrix}$ is a vector of \mathbb{C}^{2^n} ,

$$\langle \psi | \stackrel{\mathrm{def}}{=} \left(\, | \psi \rangle \, \right)^\dagger = \left(\overline{\alpha_1} \quad \dots \quad \overline{\alpha_{2^n}} \right) \text{ is a } \mathsf{bra} \, \left(\mathsf{don't} \, \mathsf{forget} \, \mathsf{the} \, \mathsf{conjugate}, \, \overline{\alpha_i}, \, \mathsf{not} \, \alpha_i \right)$$

Useful notation:

 \longrightarrow It enables to interpret $\langle \psi | \varphi \rangle$ as $\langle \psi | \cdot | \varphi \rangle$

Bra	Ket
$\langle \psi $	$ \psi\rangle$

THE KET-BRA NOTATION:

The $|\varphi\rangle\langle\psi|$ operator:

$$\begin{split} |\varphi\rangle\langle\psi| : \left(\mathbb{C}^2\right)^{\otimes n} &\longrightarrow \left(\mathbb{C}^2\right)^{\otimes n} \\ |\psi'\rangle &\longmapsto |\varphi\rangle\langle\psi| \left|\psi'\right\rangle \stackrel{\text{def}}{=} \left\langle\psi|\psi'\right\rangle |\varphi\rangle \,. \end{split}$$

Exercise:

- 1. Give the image of $|0\rangle$ and $|1\rangle$ by $|0\rangle\langle 1| + |1\rangle\langle 0|$. Give the matrix representation of this operator. Do you recognize a quantum gate?
- 2. Let $(|i\rangle)_{i\in\mathcal{I}}$ be an orthonormal basis. Which operator is

$$\sum_{i \in \mathcal{T}_i} |i\rangle\langle i|?$$

ADJOINT OF AN OPERATOR:

Adjoint of an operator:

 \mathbf{A}^{\dagger} is known as the adjoint of \mathbf{A}

Exercise:

- 1. Show that $\left(\mathbf{A}\left|\varphi\right\rangle\right)^{\dagger}=\left\langle \varphi\right|\mathbf{A}^{\dagger}$
- 2. Show that $\Big(\,|\varphi\rangle\!\langle\psi|\,\Big)^{\dagger}=|\psi\rangle\!\langle\varphi|$

Be careful with adjoint/dagger over tensor product! (do not reverse the order...)

Proposition:

We have.

$$\Big(\left.\left|\varphi\right\rangle\left.\left|\psi\right\rangle\right.\Big)^{\dagger} = \left.\left\langle\varphi\right|\left.\left\langle\psi\right|\right., \quad \left(\mathbf{A}\otimes\mathbf{B}\right)^{\dagger} = \mathbf{A}^{\dagger}\otimes\mathbf{B}^{\dagger} \quad \text{and} \quad \left(\mathbf{U}\left(\left.\left|\varphi\right\rangle\left.\left|\psi\right\rangle\right.\right)\right.\Big) = \left.\left\langle\varphi\right|\left.\left\langle\psi\right|\mathbf{U}^{\dagger}\right\rangle\right.$$

Proof:

Use the definition of tensor product as multiplication raw/column.

Classical error:

$$\Big(\left.\left|\varphi\right\rangle\left.\left|\psi\right\rangle\right.\Big)^{\dagger} = \left\langle\psi\right|\left\langle\varphi\right| \quad \text{and} \quad \Big(\mathbf{A}\otimes\mathbf{B}\Big)^{\dagger} = \mathbf{B}^{\dagger}\otimes\mathbf{A}^{\dagger}$$

