

# LECTURE 1

## INTRODUCTION TO QUANTUM COMPUTING

INF587 Quantum computer science and applications

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## Feynman (1981)

Can quantum systems be probabilistically simulated by a classical computer?

→ The answer is almost certainly, no!

→ Use quantum systems/computers to simulate quantum systems!

(birth of quantum simulation)

## A natural question

What other problems can quantum computers solve more quickly than classical computer?

## Deutsch (1985)

Foundation of quantum computing!

→ Deutsch-Jozsa algorithm (1992) quantum algorithm faster than any classical algorithm

### Shor (1994)

Solves the discrete logarithm and factoring problem in polynomial time with a quantum computer!

**Terrible situation:** public-key cryptography currently deployed is broken by using an “efficient” quantum computer

→ Crypto community worried about this since many years...

There exists quantum resistant solutions: post-quantum cryptography (active research topic)

American's government (2017) has launched a process to standardized post-quantum primitives

### Grover (1996)

Find an element in a list of size  $n$  in time  $O(\sqrt{n})$  while any classical algorithm needs a time  $\approx n$

**Consequence:** size of keys in symmetric cryptography has to be  $\times 2$ .

Computations are “noisy”

- ▶ Quantum bits are very fragile, they quickly interfere with the environment: **decoherence**
- ▶ Quantum architectures are not “ideal”

→ Faults in computation can theoretically be “corrected”: **quantum error correcting codes**

**Theorem 1. [Aharonov, Ben-Or, 1997]**

Quantum computation is possible provided the noise is sufficiently low

## Benett-Brassard (1984)

Quantum protocol for key-exchange

- ▶ Already implemented
- ▶ If an authenticated canal has been established, **unconditional security**: relies only on the validity of physic laws and not **computational assumptions**

→ Basics of **quantum computing** and **quantum information** theory

- Quantum formalism with density operators, general measures, partial trace, etc...
- Quantum circuit model, quantum algorithms (Deutsch-Josza, Simon, Grover, Quantum Fourier Transform, Shor)
- Basics of quantum error correcting codes and quantum cryptography

### References:

- ▶ Nielsen and Chuang, *Quantum computation and quantum information*,  
→ Nice introduction to quantum computing and quantum information
- ▶ de Wolf's lecture notes: <https://arxiv.org/abs/1907.09415>,  
→ Nice for advanced quantum algorithms
- ▶ Zemor's lecture notes: <https://www.math.u-bordeaux.fr/~gzemor/QuantumCodes.pdf>,  
→ Introduction to quantum error correcting codes

1. An exam (3 hours): **an A3 sheet allowed**

→ Three exercises seen during the Exercise Sessions will be at the exam.

2. Presentation of a research article or a chapter of some lecture notes (30min).



You are in a course of **computer science**

Computer science: **art of computing**

—→ We don't care that an object "exists", we want to compute it **efficiently!**

**Using the law of quantum physic: new model of computation**

What does mean **quantum computing?** What is a **quantum algorithm?**

—→ This course is not about the law of physics or about the "technologies" to verify/use them

# CLASSICAL BITS VERSUS QUANTUM BITS

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► Classical bit:  $b \in \{0, 1\}$

► Probabilistic bit:  $\begin{pmatrix} p \\ q \end{pmatrix}$  where

$$p \stackrel{\text{def}}{=} \mathbb{P}(b = 0)$$

$$q \stackrel{\text{def}}{=} \mathbb{P}(b = 1)$$

► Evolution during a computation (a probabilistic bit stays a probabilistic bit):

$$\begin{pmatrix} p \\ q \end{pmatrix} \longrightarrow \begin{pmatrix} p' \\ q' \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} p \\ q \end{pmatrix} \quad \text{where } \begin{cases} a + c = 1 \\ b + d = 1 \end{cases} \quad \text{and } a, b, c, d \geq 0.$$

Probabilistic computation: multiplication by a **stochastic** matrix

Examples:  $b \rightarrow b \oplus b$  and  $b \mapsto b \oplus 1$

$$\begin{pmatrix} p \\ q \end{pmatrix} \longrightarrow \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} p \\ q \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} p \\ q \end{pmatrix} \longrightarrow \begin{pmatrix} q \\ p \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} p \\ q \end{pmatrix}$$

*"A superposition of classical states"*

- ▶ A qubit  $|\psi\rangle$  is an element of  $\mathbb{C}^2$  with Euclidean norm 1:

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle \text{ with } \alpha, \beta \in \mathbb{C} \text{ (called amplitude) and } |\alpha|^2 + |\beta|^2 = 1$$

where  $(|0\rangle, |1\rangle)$  orthonormal basis of  $\mathbb{C}^2$ . Usually defined as

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \text{and} \quad |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

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We "cannot see" a superposition, we "can only see" classical states: measure and observe!

- ▶ **Measurement:** probabilistic orthogonal projection. Given  $|e_0\rangle, |e_1\rangle \in \mathbb{C}^2$  orthonormal basis:

$$\text{Measuring in the basis } (|e_0\rangle, |e_1\rangle) : |\psi\rangle = \alpha |e_0\rangle + \beta |e_1\rangle \xrightarrow{\text{measure}} \begin{cases} |e_0\rangle \text{ with prob. } |\alpha|^2 \\ |e_1\rangle \text{ with prob. } |\beta|^2 \end{cases}$$

## Exercise: Computational versus Hadamard basis

1. Show that  $(|+\rangle, |-\rangle)$  is an orthonormal basis of  $\mathbb{C}^2$  where

$$|+\rangle \stackrel{\text{def}}{=} \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \quad \text{and} \quad |-\rangle \stackrel{\text{def}}{=} \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

2. Give the outcome distribution when measuring  $|0\rangle, |-\rangle$ , and  $\frac{1}{\sqrt{3}} |0\rangle + \sqrt{\frac{2}{3}} |1\rangle$  in the bases  $(|0\rangle, |1\rangle)$  and  $(|+\rangle, |-\rangle)$ .

- ▶ **Qubit:**  $|\psi\rangle \in \mathbb{C}^2$  of Hermitian norm 1,
- ▶ **Measuring** in the orthonormal basis  $(|e_0\rangle, |e_1\rangle)$ :

$$|\psi\rangle = \alpha |e_0\rangle + \beta |e_1\rangle \xrightarrow{\text{measure}} \begin{cases} |e_0\rangle & \text{with prob. } |\alpha|^2 \\ |e_1\rangle & \text{with prob. } |\beta|^2 \end{cases}$$

A measurement is a “computation” you have access to

→ See Lecture 2 for a precise definition of measurement..

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Are there other computations over qubits we have access to?

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Are there other computations over qubits we have access to?

→ Yes! **Unitary evolutions**



- ▶ Unitary evolution:  $U \in \mathbb{C}^{2 \times 2}$  **unitary matrix**  $\iff UU^\dagger = I_2$

$$|\psi\rangle \longrightarrow U|\psi\rangle$$

Is it true that a qubit is still a qubit after a unitary evolution? Why?

→ Yes! Unitary evolutions preserve the Hermitian norm (more generally the inner-product)

**Unitary evolutions are invertible!**

$$|\psi\rangle \xrightarrow{U} U|\psi\rangle \xrightarrow{U^\dagger} U^\dagger U|\psi\rangle = |\psi\rangle$$

- ▶  $U \in \mathbb{C}^{2 \times 2}$  unitary over qubits is often called **quantum gate**

→ It exists a small set of gates which is **universal**

To define a quantum gate: enough to specify the image of an **orthonormal basis** and then extended it by linearity

**But** it has to map an orthonormal basis to an orthonormal basis!

## Exercise: Quantum Gates?

Are the following linear operators over qubits be quantum gates?

- $|0\rangle \mapsto |1\rangle$  and  $|1\rangle \mapsto \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$ ,
- $|0\rangle \mapsto |1\rangle$  and  $|1\rangle \mapsto |0\rangle$ .

Quantum gates have matrix representations!

For instance:  $|0\rangle \mapsto |1\rangle$  and  $|1\rangle \mapsto |0\rangle$  has the representation:  $\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ . Only linear operator that maps  $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$  to  $|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$  and  $|1\rangle$  to  $|0\rangle$ .

## ▶ NOT-gate X:

Linear op.	Matrix rep.
$ 0\rangle \mapsto  1\rangle$ $ 1\rangle \mapsto  0\rangle$	$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

## ▶ Hadamard-gate H:

Linear op.	Matrix rep.
$ 0\rangle \mapsto \frac{1}{\sqrt{2}} ( 0\rangle +  1\rangle)$ $ 1\rangle \mapsto \frac{1}{\sqrt{2}} ( 0\rangle -  1\rangle)$	$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$

## Exercise:

1. What is the effect of applying H on  $|0\rangle$  and measuring it?
2. What is the effect of applying H on  $|0\rangle$  twice?

Is quantum computation over qubits the same than classical computation over probabilistic bits?

## Exercise:

Show that there is no stochastic matrix  $\mathbf{P}$  which when applied to  $0$ , *i.e.* to  $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$ , simulates the effect of the Hadamard gate

The “ $-1$ ” gives you a huge power...

# YOUR FIRST QUANTUM ALGORITHM

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## Problem

- Input:  $f : \{0, 1\}^n \rightarrow \{0, 1\}$  either constant or balanced,
- Output: 0 if and only if  $f$  is constant.

Query complexity to  $f$ :

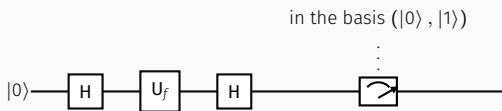
- ▶ Classically:  $1 + \frac{2^n}{2}$ ,
- ▶ Quantumly: 1.

# THE DEUTSCH-JOSZA ALGORITHM FOR $n = 1$

- ▶ Suppose that we have access to the following gate (see exercise session)

$$|b\rangle \longrightarrow \boxed{U_f} \longrightarrow (-1)^{f(b)} |b\rangle$$

- ▶ The algorithm



- ▶ Analysis

1. Applying  $H$ :  $\frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$ ,
2. Applying  $U_f$ :

$$U_f \left( \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \right) = \frac{1}{\sqrt{2}} (U_f |0\rangle + U_f |1\rangle) = \frac{(-1)^{f(0)} |0\rangle + (-1)^{f(1)} |1\rangle}{\sqrt{2}}$$

3. Applying  $H$ :

$$\begin{aligned} H \left( \frac{(-1)^{f(0)} |0\rangle + (-1)^{f(1)} |1\rangle}{\sqrt{2}} \right) &= \frac{1}{\sqrt{2}} \left( (-1)^{f(0)} H |0\rangle + (-1)^{f(1)} H |1\rangle \right) \\ &= \frac{\left( (-1)^{f(0)} + (-1)^{f(1)} \right) |0\rangle + \left( (-1)^{f(0)} - (-1)^{f(1)} \right) |1\rangle}{2} \end{aligned}$$

Before measuring we have computed

$$|\psi_{\text{out}}\rangle \stackrel{\text{def}}{=} \frac{\left((-1)^{f(0)} + (-1)^{f(1)}\right) |0\rangle + \left((-1)^{f(0)} - (-1)^{f(1)}\right) |1\rangle}{2}$$

► If  $f$  constant:

$$|\psi_{\text{out}}\rangle = \pm |0\rangle .$$

► If  $f$  balanced, namely  $f(0) \neq f(1)$ :

$$|\psi_{\text{out}}\rangle = \pm |1\rangle .$$



Before measuring we have computed

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► If  $f$  balanced, namely  $f(0) \neq f(1)$ :

$$|\psi_{\text{out}}\rangle = \pm |1\rangle .$$

Measuring in the  $(|0\rangle, |1\rangle)$  basis leads to (with **probability one**)

$|0\rangle$  if  $f$  constant or  $|1\rangle$  if  $f$  balanced

# $n$ QUBITS SYSTEM

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During all this course we will work in **finite dimension**, think  $\mathbb{C}^N$

→ Vector spaces have finite dimension, linear operator can be written as matrices, etc...

Given two vector spaces  $V$  and  $W$ , the **tensor product**  $\mathbf{v} \otimes \mathbf{w}$  between  $\mathbf{v} \in V$  and  $\mathbf{w} \in W$  verifies:

(1) for any scalar  $z$ ,

$$z(\mathbf{v} \otimes \mathbf{w}) = (z\mathbf{v}) \otimes \mathbf{w} = \mathbf{v} \otimes (z\mathbf{w})$$

(2) for any  $\mathbf{v}_1, \mathbf{v}_2 \in V$ ,

$$(\mathbf{v}_1 + \mathbf{v}_2) \otimes \mathbf{w} = \mathbf{v}_1 \otimes \mathbf{w} + \mathbf{v}_2 \otimes \mathbf{w}$$

(3) for any  $\mathbf{w}_1, \mathbf{w}_2 \in W$ ,

$$\mathbf{v} \otimes (\mathbf{w}_1 + \mathbf{w}_2) = \mathbf{v} \otimes \mathbf{w}_1 + \mathbf{v} \otimes \mathbf{w}_2$$

Think the tensor product  $\mathbf{v} \otimes \mathbf{w}$  as a column/row product:

$$\begin{pmatrix} v_1 \\ \vdots \\ v_N \end{pmatrix} \begin{pmatrix} w_1 & \cdots & w_{N'} \end{pmatrix}$$

## Tensor Product of Spaces

$V$  and  $W$  be two vector spaces with bases the  $\mathbf{v}_i$ 's and the  $\mathbf{w}_j$  respectively

$$V = \text{Span}(\mathbf{v}_1, \dots, \mathbf{v}_n) \quad \text{and} \quad W = \text{Span}(\mathbf{w}_1, \dots, \mathbf{w}_m)$$

The vector space  $V \otimes W$  is defined as being generated by the  $\mathbf{v}_i$ 's and the  $\mathbf{w}_j$ 's

$$V \otimes W \stackrel{\text{def}}{=} \text{Span}(\mathbf{v}_i \otimes \mathbf{w}_j : 1 \leq i \leq n, 1 \leq j \leq m)$$

► **Dimension,**

$$\dim V \otimes W = \dim V \dim W = nm$$

► **Basis,**  $(\mathbf{v}_1, \dots, \mathbf{v}_n)$  (resp.  $(\mathbf{w}_1, \dots, \mathbf{w}_m)$ ) be a basis of  $V$  (resp.  $W$ )

$(\mathbf{v}_i \otimes \mathbf{w}_j : 1 \leq i \leq n, 1 \leq j \leq m)$  is a basis of  $V \otimes W$

► **Characterization,**

$$\mathbf{x} \in V \otimes W \iff \exists \alpha_{i,j} : \mathbf{x} = \sum_{\substack{1 \leq i \leq n \\ 1 \leq j \leq m}} \alpha_{i,j} \mathbf{v}_i \otimes \mathbf{w}_j$$

## Classical Error:

$\mathbf{x} \in V \otimes W$ , then there exists  $\mathbf{v} \in V$  and  $\mathbf{w} \in W$  such that  $\mathbf{x} = \mathbf{v} \otimes \mathbf{w}$ .

$(\mathbf{v}_1, \dots, \mathbf{v}_n)$  ( resp.  $(\mathbf{w}_1, \dots, \mathbf{w}_m)$  ) be a basis of  $V$  ( resp.  $W$  ).

## Scalar product over tensor product spaces

Suppose that  $V$  ( resp.  $W$  ) is equipped by a scalar product  $\langle \cdot, \cdot \rangle_V$  ( resp.  $\langle \cdot, \cdot \rangle_W$  ). The scalar product over  $V \otimes W$  is defined as (and extended by bilinearity) as

$$\langle \mathbf{v}_i \otimes \mathbf{w}_j, \mathbf{v}_k \otimes \mathbf{w}_\ell \rangle_{V \otimes W} \stackrel{\text{def}}{=} \langle \mathbf{v}_i, \mathbf{v}_k \rangle_V \langle \mathbf{w}_j, \mathbf{w}_\ell \rangle_W$$

## An important remark

If  $\mathbf{v}_1 \perp \mathbf{v}_2$ , then for all  $\mathbf{w}_1, \mathbf{w}_2$ :  $\mathbf{v}_1 \otimes \mathbf{w}_1 \perp \mathbf{v}_2 \otimes \mathbf{w}_2$

$(\mathbf{v}_1, \dots, \mathbf{v}_n)$  ( resp.  $(\mathbf{w}_1, \dots, \mathbf{w}_m)$  ) be a basis of  $V$  (resp.  $W$ ).

## Linear Operator over tensor product of spaces

Given  $\mathbf{A}, \mathbf{B}$  be linear operator over  $V, W$ ,  $\mathbf{A} \otimes \mathbf{B}$  is a linear operator over  $V \otimes W$  be defined (and extended by linearity) as

$$\mathbf{A} \otimes \mathbf{B} (\mathbf{v}_i \otimes \mathbf{w}_j) \stackrel{\text{def}}{=} \mathbf{A}\mathbf{v}_i \otimes \mathbf{B}\mathbf{w}_j$$

### ► Characterization,

$$\mathbf{C} \text{ linear operator over } V \otimes W \iff \exists \alpha_i, \mathbf{A}_i, \mathbf{B}_i : \mathbf{C} = \sum_i \alpha_i \mathbf{A}_i \otimes \mathbf{B}_i$$

### Classical Error:

$\mathbf{C}$  linear operator over  $V \otimes W$ , then there exists  $\mathbf{A}, \mathbf{B}$  linear operators over  $V$  and  $W$  s.t  $\mathbf{C} = \mathbf{A} \otimes \mathbf{B}$ .

## Tensor product of matrices

Let  $\mathbf{A} \stackrel{\text{def}}{=} (a_{i,j})_{\substack{1 \leq i \leq n \\ 1 \leq j \leq m}} \in \mathbb{C}^{n \times m}$  and  $\mathbf{B} \in \mathbb{C}^{p \times q}$ , then

$$\mathbf{A} \otimes \mathbf{B} \stackrel{\text{def}}{=} \begin{pmatrix} a_{1,1}\mathbf{B} & a_{1,2}\mathbf{B} & \cdots & a_{1,m}\mathbf{B} \\ a_{2,1}\mathbf{B} & a_{2,2}\mathbf{B} & \cdots & a_{2,m}\mathbf{B} \\ \vdots & \vdots & \vdots & \vdots \\ a_{n,1}\mathbf{B} & a_{n,2}\mathbf{B} & \cdots & a_{n,m}\mathbf{B} \end{pmatrix} \in \mathbb{C}^{np \times mq}$$

## Example:

$$1. \begin{pmatrix} 1 \\ 2 \end{pmatrix} \otimes \begin{pmatrix} 2 \\ 3 \end{pmatrix} = \begin{pmatrix} 1 \times 2 \\ 1 \times 3 \\ 2 \times 2 \\ 2 \times 3 \end{pmatrix} = \begin{pmatrix} 2 \\ 3 \\ 4 \\ 6 \end{pmatrix}.$$

$$2. \mathbf{X} \otimes \mathbf{H} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \otimes \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 0 & 1 & 1 \\ 0 & 0 & 1 & -1 \\ 1 & 1 & 0 & 0 \\ 1 & -1 & 0 & 0 \end{pmatrix}$$



Properties: for any  $\alpha \in \mathbb{C}$ ,  $A, B \in \mathbb{C}^{m \times n}$  and  $C, D \in \mathbb{C}^{p \times q}$

1.  $\alpha (A \otimes C) = (\alpha A) \otimes C = A \otimes (\alpha C)$ ,

2.  $(A + B) \otimes C = A \otimes C + B \otimes C$ ,

3.  $C \otimes (A + B) = C \otimes A + C \otimes B$ ,

4. If we can form matrix products  $AC$  and  $BD$ , then

$$(A \otimes B) (C \otimes D) = (AC) \otimes (BD)$$

5. If  $A, B$  are invertible, then

$$(A \otimes B)^{-1} = A^{-1} \otimes B^{-1}.$$

Classical Error:

$$A \otimes B = B \otimes A.$$

## $n$ QUBITS SYSTEM

- ▶ A qubit  $|\psi\rangle$  is an element of  $\mathbb{C}^2$  with Hermitian norm 1,
- ▶ A **register of  $n$  qubits**  $|\psi\rangle$  is an element of  $\underbrace{\mathbb{C}^2 \otimes \dots \otimes \mathbb{C}^2}_{n \text{ times}} = \mathbb{C}^{2^n}$  with Euclidean norm 1.

Let  $(|0\rangle, |1\rangle)$  be an orthonormal basis of  $\mathbb{C}^2$ . Then,

$$(|b_1\rangle \otimes |b_2\rangle \otimes \dots \otimes |b_n\rangle : b_1, \dots, b_n \in \{0, 1\})$$

is an orthonormal basis of  $\underbrace{\mathbb{C}^2 \otimes \dots \otimes \mathbb{C}^2}_{n \text{ times}} = \mathbb{C}^{2^n}$ .

- ▶ Notation: for  $b_1, \dots, b_n \in \{0, 1\}$  and  $|\psi_1\rangle, |\psi_2\rangle, \dots, |\psi_n\rangle$  be qubits  
 $|b_1 b_2 \dots b_n\rangle \stackrel{\text{def}}{=} |b_1\rangle \otimes |b_2\rangle \otimes \dots \otimes |b_n\rangle$  and  $|\psi_1\rangle |\psi_2\rangle \dots |\psi_n\rangle \stackrel{\text{def}}{=} |\psi_1\rangle \otimes |\psi_2\rangle \otimes \dots \otimes |\psi_n\rangle$
- ▶ Characterization: any register  $|\psi\rangle \in \mathbb{C}^{2^n}$  of  $n$  qubits can be written as

$$|\psi\rangle = \sum_{\mathbf{x} \in \{0,1\}^n} \alpha_{\mathbf{x}} |\mathbf{x}\rangle \quad \text{where } \alpha_{\mathbf{x}} \in \mathbb{C} \text{ (called amplitude) and } \sum_{\mathbf{x} \in \{0,1\}^n} |\alpha_{\mathbf{x}}|^2 = 1.$$

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**A remark: choose your orthonormal basis!**

From any  $(|e_0\rangle, |e_1\rangle)$  orthonormal basis of  $\mathbb{C}^2$ , then  $(|e_{i_1}\rangle \dots |e_{i_n}\rangle)$  for  $i_1, \dots, i_n \in \{0, 1\}^n$  is an orthonormal basis of  $\underbrace{\mathbb{C}^2 \otimes \dots \otimes \mathbb{C}^2}_{n \text{ times}} = \mathbb{C}^{2^n}$ .

## Exercise:

1. Compute the scalar product between  $|+\rangle |1\rangle$ ,  $|00\rangle$  and  $|11\rangle$  where  $|+\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$ .
2. Let  $(|e_0\rangle, |e_1\rangle)$  be an orthonormal basis of  $\mathbb{C}^2$ . Show that  $(|e_{i_1}\rangle \dots |e_{i_n}\rangle)$  for  $i_1, \dots, i_n \in \{0, 1\}^n$  is an orthonormal basis of  $\underbrace{\mathbb{C}^2 \otimes \dots \otimes \mathbb{C}^2}_{n \text{ times}} = \mathbb{C}^{2^n}$ .
3. Do we have  $|00\rangle + |10\rangle = (|0\rangle + |1\rangle) \otimes |0\rangle$ ?
4. (\*) Do there exist two qubits  $|\psi_1\rangle$  and  $|\psi_2\rangle$  such that

$$\frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) = |\psi_1\rangle \otimes |\psi_2\rangle .$$

5. Do there exist two qubits  $|\psi_1\rangle$  and  $|\psi_2\rangle$  such that

$$\frac{1}{2} (|00\rangle + |01\rangle + |10\rangle + |11\rangle) = |\psi_1\rangle \otimes |\psi_2\rangle .$$

## Separable versus entangled states:

A  $n$ -qubit system  $|\psi\rangle$  that can be decomposed as  $|\psi\rangle = |\psi_1\rangle \otimes |\psi_2\rangle$  is called **separable**.  
When there is no such decomposition, the state is called **entangled**.

## Example:

1. Separable states

$$|00\rangle = |0\rangle \otimes |0\rangle \quad \text{and} \quad \frac{1}{2} (|00\rangle + |01\rangle + |10\rangle + |11\rangle) = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

2. Entangled state

$$\frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

→ **Entangled states play a crucial role** in quantum computation/information (teleportation, quantum cryptography, ...)

- ▶ **Measuring in the basis**  $|e_1\rangle |e_2\rangle \dots |e_n\rangle$ :

$$|\psi\rangle = \sum_{i_1, \dots, i_n \in \{0,1\}^n} \alpha_{i_1 \dots i_n} |e_{i_1}\rangle \dots |e_{i_n}\rangle \xrightarrow{\text{measure}} |e_{j_1}\rangle \dots |e_{j_n}\rangle \text{ with probability } |\alpha_{j_1 \dots j_n}|^2$$

- ▶ **Measuring the first register in the basis**  $(|e_0\rangle, |e_1\rangle)$

$$|\psi\rangle = \alpha_0 |e_0\rangle |\psi_0\rangle + \alpha_1 |e_1\rangle |\psi_1\rangle \xrightarrow{\text{measure}} \begin{cases} |e_0\rangle |\psi_0\rangle \text{ with prob. } |\alpha_0|^2 \\ |e_1\rangle |\psi_1\rangle \text{ with prob. } |\alpha_1|^2 \end{cases}$$

Be careful: necessarily  $|\alpha_0|^2 + |\alpha_1|^2 = 1$ .

## Exercise:

Give the outcome distribution of measuring in the basis  $(|bb'\rangle : b, b' \in \{0, 1\})$  the first registers of the following two-qubits

$$|0\rangle \left( \sqrt{\frac{1}{3}} |0\rangle + \sqrt{\frac{2}{3}} |1\rangle \right), \quad \sqrt{\frac{1}{2}} |01\rangle + \sqrt{\frac{1}{3}} |11\rangle + \sqrt{\frac{1}{6}} |10\rangle \quad \text{and} \quad \frac{1}{2} (|0\rangle - |1\rangle)(|0\rangle - |1\rangle)$$

Unitary evolution  $U \in \mathbb{C}^{2^n \times 2^n}$  unitary matrix  $\iff UU^\dagger = I_{2^n}$ .

**Exercise:**

Is the following operator a unitary of  $\mathbb{C}^2 \otimes \mathbb{C}^2$ :

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

Describe the image of  $|bb'\rangle$  for  $b, b' \in \{0, 1\}$ .

# BRA-KET AND KET-BRA NOTATION

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## Scalar Product:

Let  $|e_1\rangle, \dots, |e_{2n}\rangle$  be an orthonormal basis,  $|\psi\rangle \stackrel{\text{def}}{=} \sum_i \alpha_i |e_i\rangle$  and  $|\varphi\rangle \stackrel{\text{def}}{=} \sum_i \beta_i |e_i\rangle$ . Then

$$\langle\psi|\varphi\rangle \stackrel{\text{def}}{=} \sum_i \overline{\alpha_i} \beta_i.$$

- ▶ **Ket-notation:**  $|\psi\rangle$  is called a ket
- ▶ **Bra-notation:** a ket  $|\psi\rangle = \begin{pmatrix} \alpha_1 \\ \vdots \\ \alpha_{2n} \end{pmatrix}$  is a vector of  $\mathbb{C}^{2n}$ ,

$\langle\psi| \stackrel{\text{def}}{=} (|\psi\rangle)^\dagger = (\overline{\alpha_1} \quad \dots \quad \overline{\alpha_{2n}})$  is a bra

## Useful notation:

→ It enables to interpret  $\langle\psi|\varphi\rangle$  as  $\langle\psi| \cdot |\varphi\rangle$ .

Bra	Ket
$\langle\psi $	$ \psi\rangle$

## The $|\varphi\rangle\langle\psi|$ operator

$$|\varphi\rangle\langle\psi| : (\mathbb{C}^2)^{\otimes n} \rightarrow (\mathbb{C}^2)^{\otimes n}$$

$$|\psi'\rangle \mapsto |\varphi\rangle\langle\psi| |\psi'\rangle \stackrel{\text{def}}{=} \langle\psi|\psi'\rangle |\varphi\rangle.$$

## Exercise:

1. Give the image of  $|0\rangle$  and  $|1\rangle$  by  $|0\rangle\langle 1| + |1\rangle\langle 0|$ . Give the matrix representation of this operator. Do you recognize a quantum gate?
2. Let  $(|i\rangle)$  be an orthonormal basis. Which operator is

$$\sum_i |i\rangle\langle i|.$$

### Adjoint of an operator

$A^\dagger$  is known as the adjoint of  $A$

### Exercise:

1. Show that  $(A|\varphi\rangle)^\dagger = \langle\varphi|A^\dagger$ ,
2. Show that  $(|\varphi\rangle\langle\psi|)^\dagger = |\psi\rangle\langle\varphi|$ .

Be careful with adjoint/dagger over tensor product... (do not reverse the order...)

### Proposition:

We have

$$(|\varphi\rangle\langle\psi|)^{\dagger} = \langle\varphi|\langle\psi| \quad \text{and} \quad (\mathbf{A} \otimes \mathbf{B})^{\dagger} = \mathbf{A}^{\dagger} \otimes \mathbf{B}^{\dagger}$$

### Proof:

See exercise session!

# EXERCISE SESSION

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